# The Structure of Geographical Threshold Graphs

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**Abstract.** We analyze the structure of random graphs generated by the geographical threshold model. The model is a generalization of random geometric graphs. Nodes are distributed in space, and edges are assigned according to a threshold function involving the distance between nodes as well as randomly chosen node weights. We show how the degree distribution, percolation and connectivity transitions, clustering coefficient, and diameter relate to the threshold value and weight distribution. We give bounds on the threshold value guaranteeing the presence or absence of a giant component, connectivity and disconnectivity of the graph, and small diameter. Finally, we consider the clustering coefficient for nodes with a given degree l, finding that its scaling is very close to 1/l when the node weights are exponentially distributed.

## I. Introduction

Large networks such as the Internet, World Wide Web, phone call graphs, infectious disease contacts, and financial transactions have provided new challenges for modeling and analysis [Bonato 05]. For example, web graphs may have billions of nodes and edges, which implies that the analysis on these graphs, i.e., processing and extracting information on these large sets of data, is "hard" [Abello et al. 02]. Extensive theoretical and experimental research has been done in web-graph modeling. Early measurements suggested that the Internet exhibits a power-law degree distribution [Faloutsos et al. 99] and that the web graph also follows a power-law distribution in in- and out-degree of links [Kleinberg et al. 99]. Modeling approaches using random graphs have attempted to capture both

© A K Peters, Ltd. 1542-7951/08 \$0.50 per page the structure and dynamics of the web graph [Kumar et al. 00, Barabási and Albert 99, Aiello et al. 00, Bollobás et al. 01, Cooper and Frieze 01].

The study of random graphs began with the introduction of the uniform random graph model [Erdős and Rényi 59, Erdős and Rényi 60]. Since then, many other models have been proposed to better capture the structure seen in realworld networks [Bollobás 01, Durrett 06]. Some examples are random graph models with a given or expected degree sequence [Molloy and Reed 95, Chung and Lu 06], threshold graphs [Mahadev and Peled 95, Hagberg et al. 06] with edges created according to a function of node weights, and graphs with an underlying geometric structure, such as random geometric graphs [Penrose 03]. In this paper we study another recent addition to this collection of models: geographical threshold graphs (GTGs), a static model for networks that includes both geometric information and node weight information.

GTGs combine the geometric structure of random geometric graphs with node properties similar to threshold graphs. The properties of this graph ensemble have only recently begun to be studied [Masuda et al. 05, Bradonjić et al. 07, Bradonjić and Kong 07]. One motivation for analyzing this model is that many real networks need to be studied with a stochastic model that is "richer" than random geometric graphs. The GTG model has been applied, for instance, in the study of wireless ad hoc networks in systems in which the wireless nodes have different communication ranges or battery power [Bradonjić and Kong 07]. In that case, the weights represent available power or bandwidth of a wireless node. By varying the weights, properties such as the diameter or degree distribution can be tuned. Other possible applications of GTGs that are yet to be explored are epidemic modeling, where the weights could represent susceptibility or infectivity of an individual; and social networks, where the weights might be related to affinity or attractiveness.

## 2. Geographical Threshold Graph Model

The GTG model is constructed from a set of n nodes placed independently and uniformly at random in a volume in  $\mathbb{R}^d$ . A nonnegative weight  $w_i$ , taken randomly and independently from a probability distribution function f(w):  $\mathbb{R}_0^+ \to \mathbb{R}_0^+$ , is assigned to each node  $v_i$  for  $i \in [n]$ . Let  $F(x) = \int_0^x f(w) dw$  be the cumulative density function. For two nodes  $v_i$  and  $v_j$  at distance r, the edge (i, j) exists if and only if the following connectivity relation is satisfied:

$$G(w_i, w_j)h(r) \ge \theta_n, \tag{2.1}$$

where  $\theta_n$  is a given threshold parameter that depends on the size of the network. The function h(r) specifies the connection probability as a function of distance and is assumed to be decreasing in r. In the following we take  $h(r) = r^{-\beta}$ , for some positive  $\beta$ , which is typical, for example, of the path-loss model in wireless networks [Bradonjić and Kong 07]. The interaction strength between nodes  $G(w_i, w_j)$  is typically taken to be symmetric (to produce an undirected graph) and either multiplicatively or additively separable, i.e., in the form of  $G(w_i, w_j) = g(w_i)g(w_j)$  or  $G(w_i, w_j) = g(w_i) + g(w_j)$ .

Some basic results have already been established. For the case of uniformly distributed nodes over a unit space it has been shown [Masuda et al. 05, Bradonjić and Kong 07] that the expected degree of a node with weight w is

$$\mathbf{E}[k(w)] = \frac{n\pi^{d/2}}{\Gamma(d/2+1)} \int_{w'} f(w') \left(h^{-1}(\theta_n/G(w,w'))\right)^d dw',$$

where  $h^{-1}$  is the inverse of h. The degree distribution has been studied for specific weight-distribution functions f(w) [Masuda et al. 05]. In both the multiplicative and additive cases of G(w, w'), questions of diameter, connectivity, and topology control have been addressed [Bradonjić and Kong 07].

Here we restrict ourselves to the case of g(w) = w,  $\beta = 2$ , and nodes distributed uniformly over a two-dimensional space. For analytical simplicity we take the space to be a unit torus, and use the additive model for the connectivity relation

$$\frac{w_i + w_j}{r^2} \ge \theta_n. \tag{2.2}$$

Certain of our techniques may be generalized to other cases in a straightforward manner.

Finally, we impose the following relatively weak conditions on the weight distribution f(w): (1) a finite mean  $\mu = E[w]$ , (2) a finite variance  $\sigma^2 = E[w^2] - E[w]^2$ , and (3) nonvanishing support everywhere on  $w \in (0, \infty)$ .

Some examples of GTG instances with exponential weight distribution  $f(w) = e^{-w}$  are shown in Figure 1.

The paper is organized as follows. We first state a basic property concerning the degree distribution of GTGs. In Section 4, Theorems 4.2 and 4.3 provide bounds on  $\theta_n$  for the absence and the existence of a giant component. Similarly, in Section 5, Theorems 5.2 and 5.3 provide bounds on  $\theta_n$  for the graph being disconnected and being connected. Section 6 gives upper bounds on the diameter, along with simulation results. Finally, in Section 7 we calculate the clustering coefficient, and discuss certain of its properties.



**Figure 1**. Instances of GTG with exponential weight distribution for n = 300 at decreasing threshold parameter values (increasing mean degree): (a)  $\theta_n/n = 2\pi$ , well below the percolation transition; (b)  $\theta_n/n = 1$ , above the percolation but below the connectivity transition; (c)  $\theta_n/n = 1/2e$ , well above connectivity.

## 3. Degree Distribution

We start by stating the degree distribution in our GTG model. Let the position vector of the nodes be  $\mathbf{x}$  and the weight vector be  $\mathbf{w}$ . Without loss of generality, let us consider node  $v_1$ . It is straightforward to show that the probability of  $v_1$  having degree k, given weights  $\mathbf{w}$ , is

$$\Pr[d_1 = k | \mathbf{w}] = \binom{n-1}{k} \prod_{i=2}^{k+1} \operatorname{area}(B(x_i, r_{i1})) \prod_{j=k+2}^n (1 - \operatorname{area}(B(x_j, r_{j1}))),$$

where  $\operatorname{area}(B(x_i, r_{i1}))$  is the area of the ball with center  $x_i$  and radius  $r_{i1}$ , and due to (2.2), the radii are given by

$$r_{i1} = \sqrt{\frac{w_1 + w_i}{\theta_n}}$$

for i = 2, ..., n. After marginalization, it follows that

$$\Pr[d_1 = k | w_1]$$

$$= \left(\prod_{i=2}^n \int_{w_i} f(w_i) dw_i\right) \Pr[d_1 = k | \mathbf{w}]$$

$$= \binom{n-1}{k} \left(\int_w f(w) \frac{\pi(w_1 + w)}{\theta_n} dw\right)^k \left(1 - \int_w f(w) \frac{\pi(w_1 + w)}{\theta_n} dw\right)^{n-1-k}$$

$$= \binom{n-1}{k} \left(\frac{\pi(w_1 + \mu)}{\theta_n}\right)^k \left(1 - \frac{\pi(w_1 + \mu)}{\theta_n}\right)^{n-1-k}.$$
(3.1)

That is, the degree distribution  $d_i$  of a node  $v_i$  with weight  $w_i$  follows the binomial distribution

$$d_i(\cdot|w_i) \sim \operatorname{Bin}(n-1, p_i), \tag{3.2}$$

where

$$p_i = \frac{\pi}{\theta_n} (w_i + \mu). \tag{3.3}$$

# 4. Giant Component

**Definition 4.1. (Giant Component.)** The *giant component* is a connected component of size  $\Theta(n)$ .

In this section we analyze the conditions for the existence of the giant component, giving bounds on the threshold parameter value  $\theta_n$  where it first appears. For  $\theta_n = gn$ , we specify positive constants g' > g'' and prove that whp (with high probability, i.e., with probability 1 - o(1)), if g > g', the giant component does not exist, whereas if g < g'', the giant component exists.

We do not prove that there is a zero-one law governing the emergence of the giant component. However, given that the probability of a giant component is zero above g' and one below g'', it would be rather surprising if the transition were not a sharp one.

#### 4.1. Absence of Giant Component

**Theorem 4.2.** Let  $\theta_n = gn$  for g > g', where  $g' = 2\pi\mu$ . Then why there is no giant component in GTG.

**Proof.** We start by introducing a slightly different GTG model from our usual one. Consider the space of possible node positions and weights,  $S = \{(x, y, w) : x, y \in [0, 1], w \ge 0\}$ . In the model we have already defined, we place n nodes in S, leading to a binomial degree distribution: call this the *binomial GTG*. Let us now instead place nodes in S according to a spatial Poisson process with rate nf(w), so that the expected number of nodes is n: call this the *Poisson GTG*. We will prove that the Poisson GTG does not have a giant component. It is straightforward to see that if the binomial GTG had a giant component with nonvanishing probability, the analogous Poisson GTG would have one as well. Thus, the binomial GTG cannot have a giant component either.

The proof's approach is similar to one given in [Alon and Spencer 00]. Divide the nodes into three classes: alive, dead, and neutral. Denote the number of alive nodes by  $Y_i$ . Now apply the following algorithm. At time t = 0, designate one node (picked uniformly at random) as being alive and all others as neutral. At each subsequent time step t, pick a node  $v_t$  uniformly at random from among those that are alive, and then consider all neutral nodes connected to  $v_t$ . Denote the number of these nodes by  $Z_t$ . Change these nodes from neutral to alive, and change  $v_t$  itself from alive to dead. The random variables  $Y_i, Z_i$  satisfy the following recurrence relation:  $Y_0 = 1$  and  $Y_t = Y_{t-1} + Z_t - 1$ , for  $t \ge 1$ . The number of alive nodes satisfies

$$Y_t - 1 = \sum_{i=1}^t Z_i - t.$$

Since neutral nodes are by definition those that have not yet been explored by the algorithm, the  $Z_i$  are independent random variables. We formalize this argument as follows. For a node  $v_i = (x_i, y_i, w_i)$ , define  $S_i \subseteq S$  as the subspace of all positions and weights of nodes that can be connected to  $v_i$ , namely  $S_i = \{(x, y, w) : x, y \in [0, 1], w \ge 0, (x - x_i)^2 + (y - y_i)^2 \le (w + w_i)/\theta_n\}$ . At time t = 0, any node within  $S_0$  is a neutral node connected to  $v_0$ . But at a subsequent time step t, nodes within any  $S_i$  for i < t have already been designated alive, so only those in

$$B_t = S_t \setminus \bigcup_{i=0}^{t-1} S_i$$

can be neutral nodes connected to  $v_t$ . Thus, the nodes figuring within  $Z_t$  and  $Z_{t'}$ , for any two different time steps t and t', are drawn from *disjoint* subspaces  $B_t$  and  $B_{t'}$ .

The  $Z_i$  are simply restrictions of the Poisson process to the  $B_i$ . Due to the memoryless nature of the Poisson process, they are independent Poisson random variables. Given that  $B_i \subseteq S_i$  and the expected population of  $S_i$  is  $np_i$  with  $p_i$  as defined in (3.3),  $Z_i$  satisfies the stochastic bound

$$\Pr[Z_i \ge k] \le \Pr[\operatorname{Po}(np_i) \ge k].$$

Now consider nodes that are alive, and let T be the largest t such that  $Y_t > 0$ . Then T is the size of the component containing  $v_0$ , and the giant component exists if and only if  $T = \Theta(n)$  with some nonvanishing probability. The variable T satisfies

$$\Pr[T \ge t] = \Pr[Y_t > 0] = \Pr[Y_t \ge 1] = \Pr\Big[\sum_{i=1}^t Z_i \ge t\Big] \le \Pr\Big[\sum_{i=1}^t \Pr(np_i) \ge t\Big].$$

We take the threshold to be  $\theta_n = gn$ . Since the sum of independent Poisson random variables is itself Poisson distributed, we need to prove that

$$\Pr\left[\operatorname{Po}\left(n\sum_{i=1}^{t}p_{i}\right)\geq t\right]\to 0$$

for  $t = \Theta(n)$ , for some g > 0. For any constant  $\epsilon \in (0, 1)$ , the following inequality holds:

$$\Pr\left[\operatorname{Po}\left(n\sum_{i=1}^{t}p_{i}\right)\geq t\right]\leq \Pr\left[\operatorname{Po}\left(n\sum_{i=1}^{t}p_{i}\right)\geq t\middle|\sum_{i=1}^{t}w_{i}\in(1\pm\epsilon)t\mu\right]+\Pr\left[\sum_{i=1}^{t}w_{i}\notin(1\pm\epsilon)t\mu\right].$$
(4.1)

We will bound the first term on the right-hand side using the concentration of Poisson random variables [Penrose 03]. To maximize the conditional probability, set  $\sum w_i = (1 + \epsilon)t\mu$ , and then let  $\lambda = n \sum_{i=1}^t p_i = (2 + \epsilon)at\mu$ , where  $a = n\pi/\theta_n = \pi/g$ . Now, given any constant  $\gamma \in (0, 1)$ , for  $t \to \infty$ , i.e.,  $\lambda \to +\infty$ , it follows that

$$\Pr[\operatorname{Po}(\lambda) \notin (1 \pm \gamma)\lambda] \le e^{-\lambda H(1-\gamma)} + e^{-\lambda H(1+\gamma)} \to 0,$$

where the function H is given by  $H(x) = 1 - x + x \log x$ , for x > 0. It is now sufficient to choose g large enough that  $t > (1 + \gamma)\lambda$ , which occurs when  $g > (2 + \epsilon)(1 + \gamma)\pi\mu$ . It follows that for any  $g > 2\pi\mu$ ,  $\epsilon$  and  $\gamma$  can be set so that the first term on the right-hand side in (4.1) goes to zero.

Now consider the second term on the right-hand side. By the central limit theorem,  $(\sum w_i - t\mu)/(\sqrt{t\sigma})$  tends to the normal distribution N(0,1) as  $t \to \infty$ , so

$$\Pr\left[\sum_{i=1}^{t} w_i \notin (1 \pm \epsilon) t \mu\right] = \Pr\left[\frac{\sum w_i - t\mu}{\sqrt{t\sigma}} \notin (-\epsilon, \epsilon) \sqrt{t} \frac{\mu}{\sigma}\right] \to 0$$

for any constant  $\epsilon$ .

Thus, for  $g > 2\pi\mu$ , the probability that T equals  $\Theta(n)$  goes to zero, and so there is no giant component.

#### 4.2. Existence of Giant Component

**Theorem 4.3.** Let  $\theta_n = gn$  for  $g < g'' = \sup_{\alpha \in (0,1)} \alpha F^{-1}(1-\alpha)/\lambda_c$ , where  $\pi \lambda_c \approx 4.52$  is the mean degree at which the giant component first appears in random geometric graphs (RGG) [Penrose 03]. Then why the giant component exists in GTG.

**Proof.** For any constant  $\alpha \in (0, 1)$ , we prove that whp there are  $\alpha n$  "highweighted" nodes, all with weights greater than or equal to some  $s_n$ ; we state  $s_n$  later. Let  $X_i$  be the indicator of the event  $w_i \geq s_n$ . Then  $\Pr[X_i = 1] = 1 - F(s_n) =: q$ . Let  $X = \sum_{i=1}^n X_i$  be the number of high-weighted nodes. Using the Chernoff bound  $\Pr[X \leq (1 - \delta) \mathbb{E}[X]] \leq \exp(-\mathbb{E}[X]\delta^2/2)$ , with  $\delta = 1 - \alpha/q$ , we obtain

$$\Pr[X \le \alpha n] = \Pr[X \le (1 - \delta) \mathbb{E}[X]] \le \exp\left(-n(q - \alpha)^2/(2q)\right) = n^{-\gamma}$$

for some constant  $\gamma > 1$  satisfying  $(q - \alpha)^2 = 2q\gamma \log n/n$ .

Solving this quadratic equation in q gives  $q = \alpha + \Theta(\log n/n)$ , so  $F(s_n) = 1 - q = 1 - \alpha - \Theta(\log n/n)$ . For any  $\epsilon > 0$  and n sufficiently large, the following inequality is satisfied:

$$F^{-1}(1-\alpha) \ge s_n \ge F^{-1}(1-\alpha-\epsilon).$$

Thus, let  $s_n$  be a sequence such that its limit satisfies

$$s_n \to F^{-1}(1-\alpha) = \Theta(1).$$

Now we consider the set of  $\alpha n$  high-weighted nodes. For each such node  $v_i$  with weight  $w_i$ , define its characteristic radius to be

$$r_t^2(w_i) = \frac{w_i}{\theta_n}.$$

Then it follows that any other high-weighted node  $v_j$  within this radius is connected to  $v_i$ , since the connectivity relation is satisfied:

$$\frac{w_i + w_j}{r^2} \ge \frac{w_i}{r_t^2} = \theta_n.$$

Let  $\theta_n = gn$ , where  $g < \alpha F^{-1}(1-\alpha)/\lambda_c$ . For the radius  $r_t$ , whp it follows that

$$r_t^2(w_i) = \frac{w_i}{\theta_n} \ge \frac{s_n}{\theta_n} > \frac{\lambda_c}{\alpha n}.$$
(4.2)

Let us therefore consider small circles with a fixed radius  $r_0$  such that  $\sqrt{s_n/\theta_n} > r_0 > \sqrt{\lambda_c/(\alpha n)}$  around each of these  $\alpha n$  nodes. A subgraph of this must be an RGG with mean degree greater than  $\pi \lambda_c$ , which whp contains a giant component. Since its size is  $\Theta(\alpha n) = \Theta(n)$ , it is a giant component of the GTG too. We may optimize the bound by taking the supremum of g over  $\alpha \in (0, 1)$ , and the theorem follows.

## 4.3. Comparison of Upper and Lower Bounds

We again stress that we have not proven a zero–one law for the emergence of the giant component. If a sharp transition does indeed exist, then g' and g'' provide bounds on its location. Here we consider the size of the gap between the two bounds.

Claim 4.4. For any weight distribution f(w),  $g'/g'' \ge 2\pi\lambda_c \approx 9.04$ .

**Proof.** First consider  $g' = 2\pi\mu$ . Using the telescope formula,  $\mu$  satisfies

$$\mu = \int_0^\infty (1 - F(y)) dy.$$

Now consider  $g'' = \sup_{\alpha \in (0,1)} \alpha F^{-1}(1-\alpha)/\lambda_c$ . We have  $F : [0, +\infty) \to [0, 1)$ . Since F is a bijection, the inverse  $F^{-1} : [0,1] \to [0, +\infty)$  exists. Let  $x = F^{-1}(1-\alpha)$ , and consequently  $\alpha = 1 - F(x)$ . Then

$$\sup_{\alpha \in (0,1)} \alpha F^{-1}(1-\alpha) = \sup_{x \in (0,\infty)} x(1-F(x)).$$

Define the function

$$h(x) = \int_0^x (1 - F(y)) dy - x(1 - F(x)).$$

Since  $h'(x) = xf(x) \ge 0$  and h(0) = 0, we know that  $h(x) \ge 0$  for every  $x \ge 0$ . Let  $x_0$  be the value at which x(1 - F(x)) has its supremum. Then,

$$\mu - \sup_{\alpha \in (0,1)} \alpha F^{-1}(1-\alpha) = \int_0^\infty (1 - F(y)) dy - x_0(1 - F(x_0)) \ge h(x_0) \ge 0,$$

from which the claim follows.

**Remark 4.5.** For the exponential distribution  $f(w) = \gamma \exp(-\gamma w)$ , we have  $g' = 2\pi/\gamma$ .

**Remark 4.6.** If  $\alpha F^{-1}(1-\alpha)$  has an extremum for  $\alpha \in (0,1)$ , this occurs at

$$\alpha = F^{-1}(1 - \alpha)f(F^{-1}(1 - \alpha)).$$

For example, for the exponential distribution the maximum is at  $\alpha = 1/e$ , giving a bound of  $g'' = 1/e\gamma\lambda_c$ .

# 5. Connectivity

**Definition 5.1.** (Connectivity.) The graph on n vertices is *connected* if the largest component has size n.

In this section we analyze conditions for connectivity, giving bounds on the threshold parameter  $\theta_n$  at which the entire graph first becomes connected. Similarly to our approach in the case of the giant component, for  $\theta_n = \kappa n/\log n$ , we specify positive constants  $\kappa' > \kappa''$  and prove that whp, if  $\kappa > \kappa'$ , the graph is disconnected, whereas if  $\kappa < \kappa''$ , the graph is connected.

As in the case of the emergence of the giant component, it seems likely but has not been proven that there is a sharp phase transition at which GTGs become connected.

#### 5.1. Disconnected Graph

In this subsection we prove the conditions under which a graph is disconnected.

**Theorem 5.2.** Let  $\theta_n = \kappa n / \log n$  for  $\kappa > \kappa'$ , where  $\kappa > \pi \mu$ . Then the GTG is disconnected whp.

**Proof.** For a node  $v_i$ , let  $Y_i$  be the indicator of the event that  $v_i$  is isolated. We will consider the sum

$$Y = \sum_{i=1}^{n} Y_i$$

and show that  $\Pr[Y = 0] \to 0$ . It will then follow that whp there is at least one isolated node, and so the graph is disconnected.

From the binomial degree distribution in (3.1), the probability that  $v_i$  is isolated, conditional on  $w_i$ , is

$$\Pr[Y_i = 1 | w_i] = \left(1 - \frac{w_i + \mu}{\theta_n} \pi\right)^{n-1}.$$

Now define

$$p \equiv \mathbf{E}[Y_i] = \Pr[Y_i = 1] = \int f(w_i) \left(1 - \frac{w_i + \mu}{\theta_n} \pi\right)^{n-1} dw_i.$$

Applying the second moment method yields

$$\Pr[Y=0] \leq \frac{\operatorname{Var}[Y]}{\operatorname{E}[Y]^2} = \frac{\sum_i \operatorname{Var}[Y_i] + \sum_{i \neq j} \operatorname{Cov}[Y_i, Y_j]}{(np)^2}$$

The variance and covariance are given by

$$Var[Y_i] = E[Y_i^2] - E[Y_i]^2 = p - p^2,$$
  

$$Cov[Y_i, Y_j] = E[Y_i, Y_j] - E[Y_i]E[Y_j] = Pr[Y_i = 1, Y_j = 1] - p^2,$$

 $\mathbf{SO}$ 

$$\Pr[Y=0] \le \frac{n(p-p^2) + n(n-1)(\Pr[Y_i=1, Y_j=1] - p^2)}{(np)^2} < \frac{1}{np} + \frac{\Pr[Y_i=1, Y_j=1]}{p^2} - 1.$$

Let us first consider the 1/(np) term. Let  $\theta_n = \kappa \frac{n}{\log n}$ , where  $\kappa$  is a constant. We claim that if  $\kappa > \pi \mu$ , then  $1/(np) \to 0$ . To see this, let  $\zeta$  be any positive constant that satisfies  $\kappa > \pi(\mu + \zeta)$ . We have

$$p = \int f(w) \left(1 - \frac{w + \mu}{\theta_n} \pi\right)^{n-1} dw$$
  

$$\geq \int_0^{\zeta} f(w) \left(1 - \frac{w + \mu}{\theta_n} \pi\right)^{n-1} dw$$
  

$$\geq F(\zeta) \left(1 - \frac{\mu + \zeta}{\theta_n} \pi\right)^{n-1}$$

$$= F(\zeta) \left(1 - \frac{\mu + \zeta}{\kappa n} \pi \log n\right)^{n-1}$$
$$= F(\zeta) n^{-(\mu + \zeta)\pi/\kappa} (1 + o(1)).$$

Therefore, if  $\kappa > \pi(\mu + \zeta)$ , then  $pn \ge F(\zeta)\omega(n)$ . Since f(w) has nonvanishing support,  $F(\zeta) > 0$ , and so  $1/(np) \to 0$  for  $n \to \infty$ .

Next, we will show that  $\Pr[Y_i = 1, Y_j = 1]/p^2 = o(1)$ . Consider the joint probability conditional on a set of weights **w**. Denoting the neighborhood relation by  $v_i \sim v_j$ , we have

$$\begin{aligned} \Pr[Y_i = 1, Y_j = 1 | \mathbf{w}] &= \Pr\left[v_i \nsim v_j, \bigcap_{k \neq i,j} v_i \nsim v_k, v_j \nsim v_k | \mathbf{w}\right] \\ &= \Pr[v_i \nsim v_j | w_i, w_j] \Pr\left[\bigcap_{k \neq i,j} v_i \nsim v_k, v_j \nsim v_k | v_i \nsim v_j, \mathbf{w}\right] \\ &= \Pr[v_i \nsim v_j | w_i, w_j] \prod_{k \neq i,j} \Pr[v_i \nsim v_k, v_j \nsim v_k | v_i \nsim v_j, \mathbf{w}]. \end{aligned}$$

We now use the easily verified property that given events Q, R, and S that depend on  $\mathbf{w}$ ,

$$\Pr[R^c, S^c | Q^c, \mathbf{w}] = 1 - \Pr[R | \mathbf{w}] - \Pr[S | \mathbf{w}] + (1 - \Pr[Q | R, S, \mathbf{w}]) \frac{\Pr[R, S | \mathbf{w}]}{\Pr[Q^c | \mathbf{w}]}$$

Let  $a = \Pr[v_i \sim v_j | w_i, w_j]$ ,  $b = \Pr[v_i \sim v_k | w_i, w_k]$ ,  $c = \Pr[v_j \sim v_k | w_j, w_k]$ , and define the *clustering coefficient* 

$$C = \Pr[v_i \sim v_j | v_i \sim v_k, v_j \sim v_k, w_i, w_j, w_k].$$

Then

$$\Pr[Y_i = 1, Y_j = 1 | \mathbf{w}] = (1 - a) \prod_{k \neq i, j} \left[ 1 - b - c + (1 - C) \frac{bc}{1 - a} \right].$$
 (5.1)

Note that  $a = (w_i + w_j)\pi/\theta_n$ , and similarly for b and c.

In Section 7.1, we show (Lemma 7.1) that if  $w_i, w_j, w_k \leq \hat{w} = (1-3\sqrt{3}/4\pi)\theta_n/2\pi$ , then  $C \geq a$ . Thus, under those conditions,

$$1 - b - c + (1 - C)\frac{bc}{1 - a} \le 1 - b - c + bc = (1 - b)(1 - c).$$

Now average (5.1) over all weights. It follows from the finite variance of f(w) that for any constant M, we have  $F(M\theta_n) = 1 - o(1/n)$ , and so

$$\begin{aligned} \Pr[Y_i &= 1, Y_j = 1] \\ &= \int f(w_i) dw_i \int f(w_j) dw_j (1-a) \\ &\times \left( \int f(w_k) dw_k \left[ 1 - b - c + (1-C) \frac{bc}{1-a} \right] \right)^{n-2} \\ &= \int_0^{\hat{w}} f(w_i) dw_i \int_0^{\hat{w}} f(w_j) dw_j (1-a) \\ &\times \left( \int_0^{\hat{w}} f(w_k) dw_k \left[ 1 - b - c + (1-C) \frac{bc}{1-a} \right] \right)^{n-2} (1+o(1)) \\ &\leq \int_0^{\hat{w}} f(w_i) dw_i \int_0^{\hat{w}} f(w_j) dw_j (1-a) \left( \int_0^{\hat{w}} f(w_k) dw_k (1-b) (1-c) \right)^{n-2} \\ &\times (1+o(1)) \\ &= \int_0^{\hat{w}} f(w_i) dw_i \int_0^{\hat{w}} f(w_j) dw_j \\ &\times \left\{ \left( 1 - \frac{\pi}{\theta_n} (w_i + w_j) \right) \\ &\times \left( \int_0^{\hat{w}} f(w_k) dw_k (1 - \frac{\pi}{\theta_n} (w_i + w_k)) (1 - \frac{\pi}{\theta_n} (w_j + w_k)) \right)^{n-2} \right\} \\ &= \int_0^{\hat{w}} f(w_i) dw_i \int_0^{\hat{w}} f(w_j) dw_j \\ &\times \left\{ (1 - \frac{\pi}{\theta_n} (w_i + w_j)) \\ &\times \left( 1 - \frac{\pi}{\theta_n} (w_i + w_j + 2\mu) + \frac{\pi^2}{\theta_n^2} (w_i w_j + \mu (w_i + w_j) + \mu^2 + \sigma^2) \right)^{n-2} \right\}. \end{aligned}$$

Now consider  $p^2$ . Using the fact that  $\mu$ ,  $\sigma$ , and  $\frac{1}{2} - (\hat{w} + \mu)\pi/\theta_n$  are all  $\Theta(1)$ , we have

$$p^{2} = \int f(w_{i})dw_{i} \int f(w_{j})dw_{j} \left(1 - \frac{\pi}{\theta_{n}}(w_{i} + \mu)\right)^{n-1} \left(1 - \frac{\pi}{\theta_{n}}(w_{j} + \mu)\right)^{n-1}$$

$$\geq \int_{0}^{\hat{w}} f(w_{i})dw_{i} \int_{0}^{\hat{w}} f(w_{j})dw_{j} \left(1 - \frac{\pi}{\theta_{n}}(w_{i} + \mu)\right)^{n-1} \left(1 - \frac{\pi}{\theta_{n}}(w_{j} + \mu)\right)^{n-1}$$

$$= \int_{0}^{\hat{w}} f(w_{i})dw_{i} \int_{0}^{\hat{w}} f(w_{j})dw_{j}$$

$$\times \left(1 - \frac{\pi}{\theta_{n}}(w_{i} + w_{j} + 2\mu) + \frac{\pi^{2}}{\theta_{n}^{2}}(w_{i}w_{j} + \mu(w_{i} + w_{j}) + \mu^{2})\right)^{n-1}$$

$$\begin{split} &= \int_{0}^{\hat{w}} f(w_{i})dw_{i} \int_{0}^{\hat{w}} f(w_{j})dw_{j} \\ &\times \left(1 - \frac{\pi}{\theta_{n}}(w_{i} + w_{j} + 2\mu) + \frac{\pi^{2}}{\theta_{n}^{2}}(w_{i}w_{j} + \mu(w_{i} + w_{j}) + \mu^{2} + \sigma^{2})\right)^{n-1}(1 - o(1)) \\ &> \int_{0}^{\hat{w}} f(w_{i})dw_{i} \int_{0}^{\hat{w}} f(w_{j})dw_{j} \left(1 - \frac{\pi}{\theta_{n}}(w_{i} + w_{j} + 2\mu)\right) \\ &\times \left(1 - \frac{\pi}{\theta_{n}}(w_{i} + w_{j} + 2\mu) + \frac{\pi^{2}}{\theta_{n}^{2}}(w_{i}w_{j} + \mu(w_{i} + w_{j}) + \mu^{2} + \sigma^{2})\right)^{n-2}(1 - o(1)) \\ &= \int_{0}^{\hat{w}} f(w_{i})dw_{i} \int_{0}^{\hat{w}} f(w_{j})dw_{j} \left(1 - \frac{\pi}{\theta_{n}}(w_{i} + w_{j})\right) \\ &\times \left(1 - \frac{\pi}{\theta_{n}}(w_{i} + w_{j} + 2\mu) + \frac{\pi^{2}}{\theta_{n}^{2}}(w_{i}w_{j} + \mu(w_{i} + w_{j}) + \mu^{2} + \sigma^{2})\right)^{n-2}(1 - o(1)) \end{split}$$

Finally, this gives the desired ratio

$$\frac{\Pr[Y_i = 1, Y_j = 1]}{p^2} < 1 + o(1).$$

By the second moment method it then follows that  $\Pr[Y=0] < o(1)$ .

#### 5.2. Connected Graphs

**Theorem 5.3.** Let  $\theta_n = \kappa n / \log n$  for  $\kappa < \sup_{\alpha \in (0,1)} \alpha F^{-1}(1-\alpha)/4$ . Then the GTG is connected whp.

**Proof.** The proof is divided into two parts. In the first part, we prove that a constant fraction of nodes  $\alpha n$  are connected. In the second part we prove that the rest of the  $(1 - \alpha)n$  nodes are connected to the first set of  $\alpha n$  nodes.

First part: Invoking the proof of the appearance of the giant component, there are  $\alpha n$  nodes all with weights greater than or equal to  $s_n \to F^{-1}(1-\alpha) = \Theta(1)$ .

Let  $\theta_n = \kappa n / \log n$ , where  $\kappa < \alpha F^{-1} (1 - \alpha) \pi$ . Analogously to  $r_t$ , define the connectivity radius  $r_c$  by

$$r_c^2(w_i) = \frac{w_i}{\theta_n} \ge \frac{s_n}{\theta_n} > \frac{\log n}{\alpha \pi n}.$$

Similarly to Theorem 4.3, let us consider small circles around each of these  $\alpha n$  nodes, and consider these nodes as an RGG. It is known that  $r_n = \sqrt{\log n/(\pi n)}$  is the connectivity threshold in RGG [Gupta and Kumar 98]. The connectivity of RGG implies the connectivity of these  $\alpha n$  nodes in our GTG.

Second part: Color the  $\alpha n$  high-weighted nodes blue, and the remaining  $(1 - \alpha)n$  nodes red. Now let us tile our space into  $n/(c_0 \log n)$  squares of size

 $c_0 \log n/n$ . We state  $c_0$  later. Consider any square  $S_i$ , and let  $B_i$  be the number of blue nodes in  $S_i$ . In expectation there are  $E[B_i] = \alpha c_0 \log n$  blue nodes in each square. From the Chernoff bound it follows that

$$\Pr[B_i \ge (1 - \delta)\alpha c_0 \log n] \ge 1 - n^{-\alpha c_0 \delta^2/2}.$$
(5.2)

Let us consider one red node r. The node r belongs to some square  $S_i$ . Let  $M_r$  be the event that the red node r is connected to some blue node  $b \in S_i$ . Let the weights of r, b be  $w_r, w_b$ , respectively. The probability of the complement of  $M_r$ , conditioned on there being at least one blue node in  $S_i$ , is given by

$$\Pr[M_r^c | B_i \ge 1] = \Pr[w_r + w_b \le r^2 \theta_n] \le \Pr\left[w_r + w_b \le 2c_0 \frac{\log n}{n} \kappa \frac{n}{\log n}\right]$$
$$= \Pr[w_r + w_b \le 2c_0 \kappa].$$

As long as  $F^{-1}(1-\alpha) > 2c_0\kappa$ , then  $w_b > 2c_0\kappa$ , and hence  $\Pr[M_r^c|B_i \ge 1] = 0$ . For large enough *n* it must hold that  $(1-\delta)\alpha c_0 \log n > 1$ , and so from (5.2), we have

$$\Pr[M_r^c] \le \Pr[M_r^c | B_i \ge (1-\delta)\alpha c_0 \log n] + \Pr[B_i < (1-\delta)\alpha c_0 \log n]$$
$$< 0 + n^{-\alpha c_0 \delta^2/2}.$$

If  $\alpha c_0 \delta^2/2 \ge 1 + \epsilon$  for some  $\epsilon > 0$ , then by the union bound,

$$\Pr[\bigcup_{r} M_{r}^{c}] \leq \sum_{r} \Pr[M_{r}^{c}] \leq (1-\alpha)nn^{-(1+\epsilon)} = (1-\alpha)n^{-\epsilon}.$$

Finally, the probability that all red nodes are connected to the set of blue nodes is given by the following relation:

$$\Pr[\bigcap_{r} M_{r}] = 1 - \Pr[\bigcup_{r} M_{r}^{c}] \ge 1 - (1 - \alpha)n^{-\epsilon} \to 1.$$
(5.3)

The requirements we have imposed on constants so far are  $\kappa < \alpha F^{-1}(1-\alpha)\pi$ ,  $\kappa < F^{-1}(1-\alpha)/(2c_0)$ , and  $\alpha c_0 \ge 2(1+\epsilon)/\delta^2$ . These conditions combine to give

$$\kappa < \alpha F^{-1}(1-\alpha) \min\left(\pi, \frac{\delta^2}{4(1+\epsilon)}\right).$$

Since  $\alpha \in (0, 1), \delta \in (0, 1)$ , and  $\epsilon > 0$  are arbitrary, we obtain

$$\kappa < \sup_{\alpha \in (0,1)} \alpha F^{-1} \frac{1-\alpha}{4}.$$



**Figure 2.** Illustration of our diameter proof technique: a sequence of adjacent squares of size  $x \times x$  links an arbitrary pair of nodes u and v in a unit-area disk.

## 6. Diameter

In this section we analyze the diameter of GTG, and provide an upper bound on it. In the design of large networks, such as the Internet and wireless networks, it is desirable to achieve low latency in the graph (i.e., the hop count between any pair of nodes in the network is small). In other words, a graph with a small diameter is desired.

We give conditions on the threshold function  $\theta_n$  such that the graph has a desired diameter in general. Furthermore, we derive conditions on  $\theta_n$ , in terms of the cumulative distribution function on weights F(w), such that diam belongs to the specific classes diam = O(1), diam =  $O(\log^q n)$ , and diam =  $O(\sqrt{n/\log n})$ . These correspond respectively to ultralow-, low-, and high-latency networks. For these three classes, we give the exact expressions on  $\theta_n$  in the case of the exponentially distributed weights. Note that all of these classes correspond to denser graphs than those we have considered so far, i.e., with  $\theta_n$  scaling as  $o(\log n/n)$  as opposed to the  $\Theta(\log n/n)$  scaling for connectivity.

Let u and v be two arbitrary nodes. Construct a sequence of adjacent squares  $S_1, S_2, \ldots, S_{O(1/x)}$ , of size  $x \times x$ , linking u and v, such that u and v are the centers of the first and last squares respectively (see Figure 2).<sup>1</sup> The geometric distance between any two nodes is  $r \leq \Theta(1)$ . Thus, there are O(1/x) squares on the straight path u - v in total.

<sup>&</sup>lt;sup>1</sup>The centers of the squares lie on the straight line u - v.

Let  $V_i$  be the number of nodes that for i = 1, 2, ..., O(1/x) lie within the square  $S_i$ . We have  $E[V_i] = nx^2$ . Using the Chernoff bound, the following is satisfied:

$$\Pr[V_i \le (1-\delta)\mathbb{E}[V_i]] \le e^{-\mathbb{E}[V_i]\delta^2/2}.$$

Taking  $\delta = 1/2$ , we get  $\Pr[V_i \leq nx^2/2] \leq e^{-nx^2/8}$ , i.e., in each square  $S_i$ , there are at least  $nx^2/2$  nodes whp.

Let  $M_i$  be the event that in a square  $S_i$ , there is at least one node with weight  $w \ge s_n$ . We will specify  $s_n$  later. We now derive a lower bound on the probability  $\Pr[M_i]$ . This probability is greater than the probability conditioned on the event that there are at least  $nx^2/2$  nodes in  $S_i$ , i.e.,

$$\Pr[M_i] \ge \Pr[M_i | V_i \ge nx^2/2] \Pr[V_i \ge nx^2/2]$$
  
$$\ge (1 - \Pr[W \le s_n]^{nx^2/2})(1 - e^{-nx^2/8})$$
  
$$= (1 - F(s_n)^{nx^2/2})(1 - e^{-nx^2/8}).$$

We now explain how we choose  $s_n$  such that any two neighboring squares  $S_j$ and  $S_{j+1}$  are connected by an edge (i.e., there are two connected nodes  $a \in S_j$ and  $b \in S_{j+1}$ ). Let weights of a and b be w and w', respectively. We showed that in any square  $S_i$  there is at least one node with weight greater than or equal to  $s_n$ , whp. We want the connectivity relation (2.1) for nodes a and b to be satisfied. The maximal distance ||a-b|| between a pair of nodes is  $||a-b|| \leq x\sqrt{5}$ . Conditioned on the events that weights w, w' are greater then  $s_n$ , we have the following relation for the connectivity of nodes a and b:

$$\Pr[a \sim b | w, w' \ge s_n] \ge \Pr[2s_n/r^2 \ge \theta_n].$$

Let us choose  $s_n = \Theta(x^2\theta_n)$ . If an arbitrary pair of nodes (u, v) is connected by a path of nodes belonging to the squares  $S_1, S_2, \ldots, S_{O(1/x)}$ , then the following relation on diam is satisfied:

$$\Pr[\text{diam} = O(1/x)] \ge \Pr[\bigcap_{i=1}^{O(1/x)} M_i] = \left( (1 - e^{-nx^2/8})(1 - F(s_n)^{nx^2/2})^{O(1/x)}, \right)$$

since the nodes, as well as weights, are distributed independently.

We now present a lemma on the diameter.

**Lemma 6.1.** Let the cumulative weight distribution function be F(w) in the GTG model. Let x and a sequence  $s_n = \Theta(x^2\theta_n)$  be such that

$$\lim_{n \to \infty} \left( 1 - F(s_n)^{nx^2/2} \right)^{1/x} = 1.$$
(6.1)

Then whp, diam = O(1/x).

**Proof.** The proof follows from the previous discussion.

#### 6.1. Some Classes of Diameter

We now analyze and state conditions on  $\theta_n$  such that diam = O(1), diam =  $O(\log^q n)$ , and diam =  $O(\sqrt{n/\log n})$ . We work out the case in which the weight distribution is exponential,  $f(w) = e^{-w}$ ,  $w \ge 0$  (i.e.,  $F(w) = 1 - e^{-w}$ ,  $w \ge 0$ ), and derive an upper bound on the threshold function  $\theta_n$  in this particular case. For some other weight distribution, the analysis would be similar. Our results are these:

- Ultralow Latency: diam = O(1). Let x < 1 be a constant and  $s_n = \theta_n$ . If  $F(\theta_n)^n \to 0$ , then diam = O(1) whp. For the exponential weight distribution it follows that  $\theta_n = o(\log n)$ .
- Low Latency: diam =  $O(\log^q n)$ . Let  $x = 1/\log^q n$  and  $s_n = \theta_n/\log^{2q} n$ . If  $F(\theta_n/\log^{2q} n)^{n/(2\log^{2q} n)}\log^q n \to 0$ , then diam =  $O(\log^q n)$  whp. For the exponential weight distribution it follows that

$$\theta_n = o\Big((\log n)^{2q(1-(\log^{2q} n)/n)}\Big).$$

• High Latency: diam =  $O(\sqrt{n/\log n})$ . Let  $x = \sqrt{\log n/n}$  and  $s_n = \theta_n \log n/n$ . If  $\sqrt{n/\log n}F(\theta_n \log n/n)^{\log n} \to 0$ , then diam =  $O(\sqrt{n/\log n})$  whp. For the exponential weight distribution it follows that

$$\theta_n = o\Big((n/\log n)^{1-1/(2\log n)}\Big).$$

We now prove these claims.

**6.1.1.** Ultralow Latency: diam = O(1). For the diameter to be a constant, let x < 1 be a constant. Invoking Lemma 6.1, it follows that diam = O(1) whp if and only if  $1 - F(s_n)^{nx^2/2} \to 1$ , i.e., if and only if  $F(s_n)^n \to 0$ . The condition on the size of diam is given by the following claim, and we can derive  $\theta_n$  such that diam = O(1) whp.

Claim 6.2. If  $F(\theta_n)^n \to 0$ , then diam = O(1) whp.

For the exponential weight distribution it follows that  $F(\theta_n)^n = (1 - e^{-\theta_n})^n \rightarrow e^{-n/e^{\theta_n}}$ . The last equation tends to zero if and only if  $n/e^{\theta_n} \rightarrow \infty$ . That is, we make the following claim.

Claim 6.3. For the exponential weight distribution  $f(w) = e^{-w}$ , the diameter in GTG satisfies diam = O(1) if  $\theta_n = o(\log n)$ .

**6.1.2.** Low Latency: diam =  $O(\log^q n)$ . Let us choose  $x = 1/\log^q n$ . Invoking Lemma 6.1, we obtain

$$(1 - F(s_n)^{nx^2/2})^{1/x} = \left(1 - F(s_n)^{n/(2\log^{2q} n)}\right)^{\log^q n}.$$
(6.2)

For  $s_n \to 0$ , the last expression tends to 1 if and only if

$$F(s_n)^{n/(2\log^{2q} n)}\log^q n \to 0,$$

where we use  $\lim_{t\to+\infty} (1-1/t)^t = 1/e$ . The condition on the size of diam is given by the following claim.

Claim 6.4. if  $F(\theta_n / \log^{2q} n)^{n/(2\log^{2q} n)} \log^q n \to 0$ , then diam =  $O(\log^q n)$  whp.

For the exponential weight distribution, the following is to be satisfied:

$$F(s_n)^{n/(2\log^{2q} n)}\log^q n = \log^q n(1 - e^{-s_n})^{n/(2\log^{2q} n)} \to s_n^{n/(2\log^{2q} n)}\log^q n \to 0,$$

or equivalently,

$$s_n = o\left((\log n)^{-\frac{2q}{n}\log^{2q}n)}\right).$$

Claim 6.5. For the exponential weight distribution  $f(w) = e^{-w}$ , the diameter in GTG satisfies diam =  $O(\log^q n)$  if  $\theta_n = o\left((\log n)^{2q(1-(\log^{2q} n)/n)}\right)$ .

**6.1.3.** High Latency: diam =  $O(\sqrt{n/\log n})$ . Let us choose  $x = \sqrt{\log n/n}$ . Invoking Lemma 6.1, we obtain

$$(1 - F(s_n)^{nx^2/2})^{1/x} = (1 - F(s_n)^{\log n})^{\sqrt{\frac{n}{\log n}}}.$$

It can be shown, using  $\lim_{t\to+\infty} (1-1/t)^t = 1/e$ , that the last expression tends to 1 if and only if  $\sqrt{n/\log n}F(s_n)^{\log n} \to 0$ . The condition on the size of diam is given by the following claim.

Claim 6.6. If  $\sqrt{n/\log n} F(\theta_n \log n/n)^{\log n} \to 0$ , then diam  $= O(\sqrt{n/\log n})$  when

For the exponential weight distribution, the following is to be satisfied:

$$\sqrt{n/\log n}F(s_n)^{\log n} = \sqrt{n/\log n}(1-e^{-s_n})^{\log n} \to \sqrt{n/\log n}s_n^{\log n} \to 0$$

or equivalently,

$$s_n = o\left( (\log n/n)^{1/(2\log n)} \right).$$



**Figure 3.** (a) For the case of diam  $= O(\log^q n)$ , with q = 1.5, the analytical solid curve is the upper bound on diam(n). Simulation results match with theoretical predictions, since the simulation points all lie below the analytical curve. (b) For the case of diam  $= O(\sqrt{n/\log n})$ , the solid curve plots the upper bound on diam(n), and this bound closely matches the experimental values.

Claim 6.7. For the exponential weight distribution  $f(w) = e^{-w}$ , the diameter in GTG satisfies

diam = 
$$O(\sqrt{n/\log n})$$
 if  $\theta_n = o\left((n/\log n)^{1-1/(2\log n)}\right)$ .

Simulation results are shown for the GTG with path-loss exponent  $\beta = 3$  (not  $\beta = 2$ ) for the case of diam =  $O(\log^{1.5} n)$  in Figure 3(a) and diam =  $O(\sqrt{n/\log n})$  in Figure 3(b). Exponentially distributed weights with mean 1 are used. The network sizes simulated are  $n = \{100, 200, 500, 1000, 2000, 10000\}$ . The threshold values  $\theta_n$  for the two cases are obtained by invoking previous claims.

# 7. Clustering Coefficient

#### 7.1. Weights Given

Let us consider in more detail the clustering coefficient defined in Section 5, namely the neighbor probability

$$C(w_i, w_j, w_k) = \Pr[v_i \sim v_j | v_i \sim v_k, v_j \sim v_k, w_i, w_j, w_k].$$

Let  $x = \sqrt{(w_i + w_j)/\theta_n}$ ,  $y = \sqrt{(w_j + w_k)/\theta_n}$ , and  $z = \sqrt{(w_i + w_k)/\theta_n}$ . Then if  $d_{ij}$  represents the distance between points *i* and *j*, we have

$$C(w_i, w_j, w_k) = \Pr[d_{ij} \le x | d_{jk} \le y, d_{ik} \le z]$$
  
=  $\frac{1}{\pi z^2} \int_0^z \Pr[d_{ij} \le x | d_{jk} \le y, d_{ik} = r] 2\pi r \, dr$   
=  $\frac{1}{\pi^2 y^2 z^2} \int_0^z A(r) 2\pi r \, dr,$ 

where A(r) is the overlap area of a disk of radius x centered at i and a disk of radius y centered at k. Now consider a triangle  $\triangle IJK$ , with sides IJ = r, IK = x, JK = y, and  $\angle JKI = \eta$ ,  $\angle IJK = \xi$ . Following arguments similar to those in [Dall and Christensen 02], there are three possible cases for A(r):

$$A(r) = \begin{cases} \pi[\min(x,y)]^2 & \text{if } r \le |x-y|, \\ x^2(\eta - \sin\eta\cos\eta) + y^2(\xi - \sin\xi\cos\xi) & \text{if } |x-y| < r < x+y, \\ 0 & \text{if } r \ge x+y, \end{cases}$$

where

$$\eta = \cos^{-1} \left( \frac{r^2 + x^2 - y^2}{2xr} \right),$$
  
$$\xi = \cos^{-1} \left( \frac{r^2 - x^2 + y^2}{2yr} \right).$$

From the definitions of x, y, and z, we have |x - y| < z < x + y. After some algebraic manipulation, one then obtains

$$C = \frac{x^2}{\pi} \left[ \frac{1}{y^2} \cos^{-1} \left( \frac{z^2 + x^2 - y^2}{2xz} \right) + \frac{1}{x^2} \cos^{-1} \left( \frac{z^2 - x^2 + y^2}{2yz} \right) \right] \\ + \frac{1}{z^2} \cos^{-1} \left( \frac{x^2 + y^2 - z^2}{2xy} \right) \\ - \frac{x^2 + y^2 + z^2}{4x^2y^2z^2} \sqrt{2x^2y^2 + 2x^2z^2 + 2y^2z^2 - x^4 - y^4 - z^4} \right] \\ = \frac{1}{\pi(w_i + w_k)(w_j + w_k)} \\ \times \left\{ (w_i + w_j)(w_i + w_k) \cos^{-1} \left( \frac{w_i}{\sqrt{w_i + w_j}\sqrt{w_i + w_k}} \right) \right. \\ + (w_i + w_j)(w_j + w_k) \cos^{-1} \left( \frac{w_j}{\sqrt{w_i + w_j}\sqrt{w_j + w_k}} \right) \\ + (w_i + w_k)(w_j + w_k) \cos^{-1} \left( \frac{w_k}{\sqrt{w_i + w_k}\sqrt{w_i + w_k}} \right) \\ - (w_i + w_j + w_k)\sqrt{w_iw_j + w_jw_k + w_kw_i} \right\}.$$

Note that while C is a function of the weights  $w_i$ ,  $w_j$ , and  $w_k$ , it is independent of  $\theta_n$ . This reflects a similar property in random geometric graphs [Dall and Christensen 02], where the clustering coefficient is independent of the graph's mean degree.

In terms of the connection probabilities a, b, and c defined in Theorem 5.2,  $x = \sqrt{a/\pi}, y = \sqrt{b/\pi}$ , and  $z = \sqrt{c/\pi}$ . Then, C is given by

$$C\pi bc = ab\cos^{-1}\frac{a+b-c}{2\sqrt{ab}} + bc\cos^{-1}\frac{-a+b+c}{2\sqrt{bc}} + ca\cos^{-1}\frac{a-b+c}{2\sqrt{ac}} - \frac{a+b+c}{4}\sqrt{2ab+2ac+2ca-a^2-b^2-c^2}.$$

We now prove the bound on the clustering coefficient that we needed for Theorem 5.2.

Lemma 7.1. If  $w_i, w_j, w_k \leq \hat{w} = (1 - 3\sqrt{3}/4\pi)\theta_n/2\pi$ , then  $C \geq a$ .

Proof. Define

$$S(a, b, c) = \frac{C\pi}{a}$$
  
=  $\frac{1}{a}\cos^{-1}\frac{-a+b+c}{2\sqrt{bc}} + \frac{1}{b}\cos^{-1}\frac{a-b+c}{2\sqrt{ac}} + \frac{1}{c}\cos^{-1}\frac{a+b-c}{2\sqrt{ab}} - \gamma\frac{a+b+c}{4abc}$ ,

where

$$\gamma = \sqrt{2ab + 2ac + 2ca - a^2 - b^2 - c^2} = 2\sqrt{bc}\sqrt{1 - \left(\frac{-a + b + c}{2\sqrt{bc}}\right)^2}$$

It is easy to verify that  $S = \pi$  when  $w_i = w_j = w_k = \hat{w}$ . We will now show that S is nonincreasing over the weights, and thus  $S \ge \pi$  for all smaller values of  $w_i, w_j, w_k$ .

Consider the sign of the derivative

$$\frac{dS}{dw_i} = \frac{\partial S}{\partial a}\frac{\partial a}{\partial w_i} + \frac{\partial S}{\partial b}\frac{\partial b}{\partial w_i} + \frac{\partial S}{\partial c}\frac{\partial c}{\partial w_i} = \frac{\pi}{\theta_n}\Big(\frac{\partial S}{\partial a} + \frac{\partial S}{\partial b}\Big).$$

Since S is symmetric in a and b, it is sufficient to consider the sign of  $\partial S/\partial a$ :

$$\frac{\partial S}{\partial a} = \frac{\gamma}{4a^2bc}(-a+b+c) - \frac{1}{a^2}\cos^{-1}\frac{-a+b+c}{2\sqrt{bc}}.$$

Now let

$$t = \frac{-a+b+c}{2\sqrt{bc}} \in [0,1].$$

Then

$$\frac{\partial S}{\partial a} = \frac{1}{a^2} \left( t \sqrt{1 - t^2} - \cos^{-1} t \right).$$

Given the function  $\varphi(t) = t\sqrt{1-t^2} - \cos^{-1} t$  on [0,1], we have  $\varphi'(t) = 2\sqrt{1-t^2} \ge 0$ , and  $\varphi(1) = 0$ . It follows that  $\partial S/\partial a \le 0$ , and so  $dS/dw_i \le 0$ . Finally, S is symmetric in  $(w_i, w_j, w_k)$ , so it must be nonincreasing over each of the weights and bounded below by the value at  $w_i = w_j = w_k = \hat{w}$ .

#### 7.2. Degree Given

Define  $C_l$  to be the neighbor probability for a node with a given degree:

$$C_{l} = \Pr[v_{i} \sim v_{j} | v_{i} \sim v_{k}, v_{j} \sim v_{k}, d(v_{k}) = l]$$

$$= \frac{\Pr[v_{i} \sim v_{j}, v_{i} \sim v_{k}, v_{j} \sim v_{k}, d(v_{k}) = l]}{\Pr[v_{i} \sim v_{k}, v_{j} \sim v_{k}, d(v_{k}) = l]}$$

$$= \frac{\int f(\mathbf{w}) \Pr[v_{i} \sim v_{j}, v_{i} \sim v_{k}, v_{j} \sim v_{k}, d(v_{k}) = l | \mathbf{w}] d\mathbf{w}}{\int f(\mathbf{w}) \Pr[v_{i} \sim v_{k}, v_{j} \sim v_{k}, d(v_{k}) = l | \mathbf{w}] d\mathbf{w}}$$

A straightforward calculation shows that  $C_l$  is given by the ratio  $I_n/I_d$  of two integrals, where

$$I_d = \int f(w_k) \left(\frac{\pi}{\theta_n}(\mu + w_k)\right)^l \left(1 - \frac{\pi}{\theta_n}(\mu + w_k)\right)^{n-l} dw_k$$

and

$$I_n = \iiint f(w_i)f(w_j)f(w_k)\pi y^2\pi z^2 C(w_i, w_j, w_k)dw_i dw_j$$
$$\times \left(\frac{\pi}{\theta_n}(\mu+w_k)\right)^{l-2} \left(1 - \frac{\pi}{\theta_n}(\mu+w_k)\right)^{n-l} dw_k$$
$$= \iiint f(w_i)f(w_j)f(w_k)(w_i+w_k)(w_j+w_k)C(w_i, w_j, w_k)dw_i dw_j$$
$$\times \frac{1}{(\mu+w_k)^2} \left(\frac{\pi}{\theta_n}(\mu+w_k)\right)^l \left(1 - \frac{\pi}{\theta_n}(\mu+w_k)\right)^{n-l} dw_k$$

For a specific weight distribution, these integrals may be evaluated numerically. It is intuitive that when l is very large  $(l = \Theta(n))$ ,  $C_l$  should scale as 1/n: nodes that connect to very many neighbors presumably do so because of their high weights, and their neighbors are no more likely to be connected than any two random nodes are. Interestingly, in the case shown in Figure 4,  $C_l$  scales almost perfectly as  $l^{-1}$  even at relatively small values of l. For an exponential weight distribution with mean 1 and parameters n = 1000 and  $\theta_n = 1000$ , the slope on the log-log plot already appears very close to -1 at l = 8.



Figure 4. Clustering coefficient vs. degree L.

## 8. Summary

Geographical threshold graphs are a rich model with the possibility of controlling structural properties by choosing specific weight distributions and tuning threshold values. The model is a versatile one and can be used for the generation and analysis not only of web graphs or large complex networks, but more generally of relation graphs in a large data set. If the data have a metric and can be mapped to nodes in Euclidean space, much of the foregoing analysis applies: one may hope to control structural properties of the data set by studying it as a GTG.

In this paper we have analyzed some of the structural properties of a GTG. Given a node weight distribution f(w) and threshold  $\theta_n$ , the degree distribution can be easily calculated.

We have given bounds on the threshold value  $\theta_n$  guaranteeing the absence or existence of the giant component. We have also given bounds on  $\theta_n$  guaranteeing a disconnected or connected graph, and provided upper bounds on the diameter for sufficiently dense graphs.

Finally, we have derived a formula for the clustering coefficient in terms of the weight distribution and threshold, as well as discussed the general clustering coefficient for nodes with a given degree, and evaluated it numerically. Our analysis has used the additive threshold function  $(w_i + w_j)/r^2 \ge \theta_n$  for the connectivity relation, but not all of our techniques require it. For this reason, many of the results may be generalized to other threshold functions, other path-loss exponents, and other spatial dimensionalities in a straightforward manner.

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